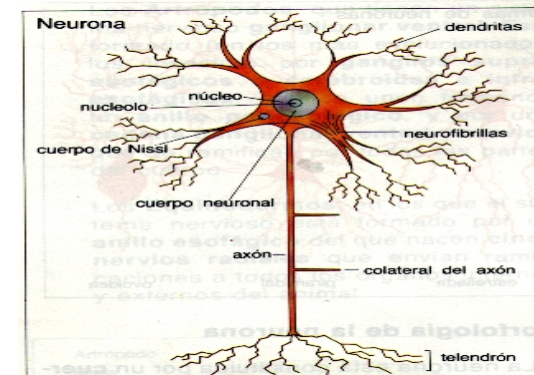
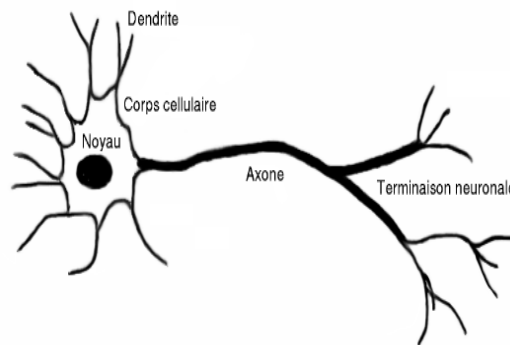
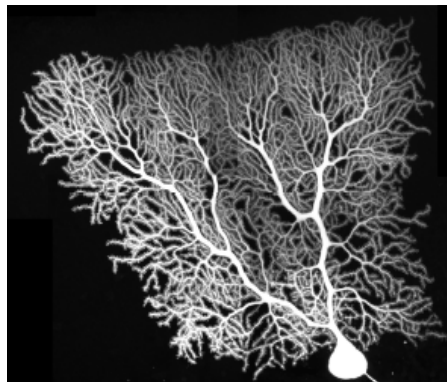


PDE models of neural networks

Benoît Perthame

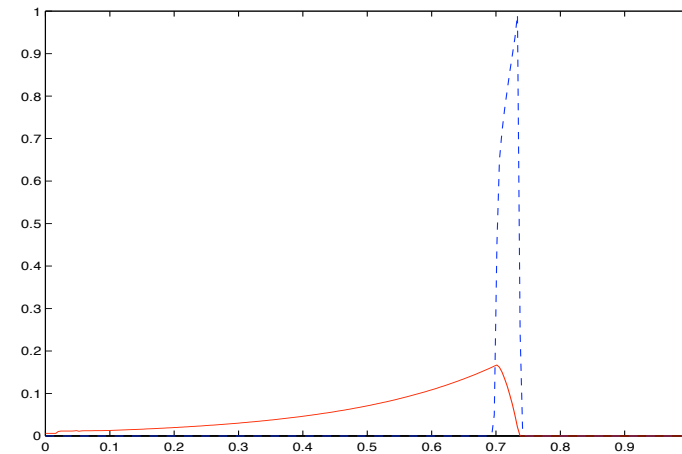
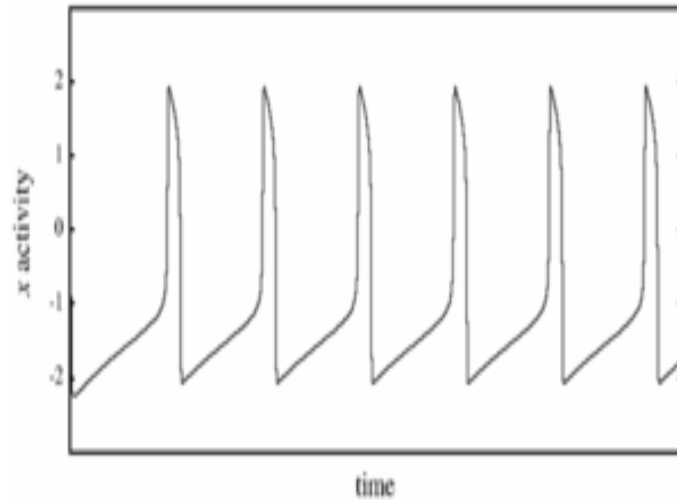


Introduction

The electrically active cells are characterized by an **action potential**

- Hodgkin-Huxley
- FitzHugh-Nagumo
- Morris-Lekar
- Izhikevich
- Mitchell-Schaeffer

Introduction



Solutions to the Hodgkin-Huxley model and to the FitzHugh-Nagumo model

These models are accurate
but very expensive/difficult to use for large assemblies of neurones.

Introduction

The **Wilson-Cowan** model (1972) describes the firing rates $N(t, x)$ of neuron assemblies located at position x through an integral equation

$$\frac{d}{dt}N(x, t) = -N(x, t) + \int w(x, y)\sigma(N(y, t))dy + s(x, t)$$

Feature : multiple steady states and bifurcation theory (Bressloff-Golubitsky, Chossat-Faugeras)

- $\sigma(\cdot)$ = sigmoid
- w_{ij} = connectivity matrix
- s = source

Introduction

The **Wilson-Cowan** model (1972) describes the firing rates $N(t, x)$ of neuron assemblies located at position x through an integral equation

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Feature : multiple steady states and bifurcation theory (Bressloff-Golubitsky, Chossat-Faugeras)

Aim : large scale brain activity, visual hallucinations (Klüver, Oster, Siegel...)



OUTLINE OF THE LECTURE

- I. Principle of Noisy Integrate and Fire model
- II. The nonlinear Noisy Integrate and Fire model
- III. The elapsed time approach

Leaky Integrate and Fire

The Leaky Integrate & Fire model is simpler

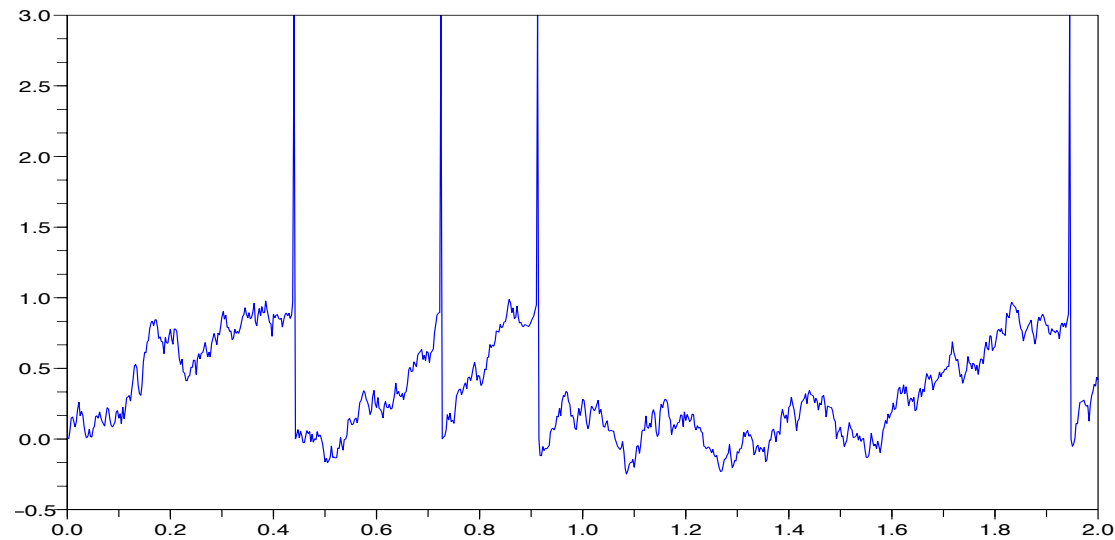
$$dV(t) = \left(-V(t) + I(t) \right) dt + \sigma dW(t), \quad V(t) < V_{\text{Firing}}$$

$$V(t_-) = V_{\text{Firing}} \implies V(t_+) = V_{\text{reset}}.$$

The idea was introduced by [L. Lapicque \(1907\)](#).

- $I(t)$ input current
- Noise or not
- Stochastic firing

Leaky Integrate and Fire



Solution to the LIF model

- N. Brunel, V. Hakim, W. Gerstner and W. Kistler...
- Fit to measurements
- Explains qualitatively observations on the brain activity

Leaky Integrate and Fire

Written in terms of PDEs, the probability $n(v, t)$ to find a neuron at the potential v

$$\left\{ \begin{array}{l} \frac{\partial n(v, t)}{\partial t} + \frac{\partial}{\partial v} \left[\overbrace{\left(-v + I(t) \right) n(v, t)}^{\text{leak+external currents}} \right] - a \overbrace{\frac{\partial^2 n(v, t)}{\partial v^2}}^{\text{Noise}} = \overbrace{\delta(v = V_R) N(t)}^{\text{neurons reset}}, \quad v \leq V_F, \\ n(V_F, t) = 0, \quad n(-\infty, t) = 0, \\ N(t) := -a \frac{\partial n(V_F, t)}{\partial v} \geq 0, \quad (\text{the total flux of neurons firing at } V_F). \end{array} \right.$$

$N(t)$ is also a Lagrange multiplier for the constraint

$$\int_{-\infty}^{V_F} n(v, t) dv = 1.$$

Leaky Integrate and Fire

$$\left\{ \begin{array}{l} \frac{\partial n(v,t)}{\partial t} + \frac{\partial}{\partial v} \left[(-v + I(t))n(v,t) \right] - a \frac{\partial^2 n(v,t)}{\partial v^2} = \delta(v = V_R)N(t), \quad v \leq V_F, \\ n(V_F, t) = 0, \quad p(-\infty, t) = 0, \\ N(t) := -a \frac{\partial n(V_F, t)}{\partial v} \geq 0, \quad (\text{the total flux of firing neurons at } V_F). \end{array} \right.$$

Properties (M. Cáceres, J. Carrillo, BP) The solutions satisfy

- $n \geq 0$, $\int_{-\infty}^{V_F} n(v, t) dv = 1$,
- For $I(t) \equiv 0$, $n(v, t) \xrightarrow[t \rightarrow \infty]{} P(v)$ the unique steady state of integral 1 (desynchronization)
- The convergence rate is exponential

Leaky Integrate and Fire

The proof uses

- the Relative Entropy

$$\frac{d}{dt} \int_{-\infty}^{V_F} P(v) H\left(\frac{n(v, t)}{P(v)}\right) dv \leq 0,$$

for $H(\cdot)$ convex,

- Hardy/Poincaré inequality,

$$\int_{-\infty}^{V_F} P(v) |u(v)|^2 dv \leq C \int_{-\infty}^{V_F} P(v) |\nabla u(v)|^2 dv,$$

for $\int_{-\infty}^{V_F} P(v) u(v) dv = 0$ [notice $P(V_F) = 0$].

Ledoux, Barthe and Roberto

Noisy LIF networks

For networks, the current $I(t)$ is related to the total activity of the network

$$\left\{ \begin{array}{l} \frac{\partial n(v,t)}{\partial t} + \frac{\partial}{\partial v} \left[(-v + bN(t))n(v,t) \right] - a(N(t)) \frac{\partial^2 n(v,t)}{\partial v^2} = \delta_{V_R}(v)N(t), \quad v \leq V_F, \\ n(V_F, t) = 0, \quad n(-\infty, t) = 0, \\ N(t) := -a(N(t)) \frac{\partial}{\partial v} n(V_F, t) \geq 0, \quad \text{total flux of firing neurons at } V_F. \end{array} \right.$$

Constitutive laws

- $I(t) = bN(t)$
- $b > 0$ for excitatory neurones
- $a(N) = a_0 + a_1N$
- $b =$ connectivity
- $b < 0$ for inhibitory neurones

Noisy LIF networks

Theorem (J. Carrillo, BP, D. Smets) Assume

- $a = a_0 > 0$ and $b < 0$ (inhibitory)
- the initial data is bounded by a supersolution (in a certain sense)

Then,

- There are global solutions
- Uniformly bounded for all $t > 0$

Noisy LIF networks

Theorem (M. Cáceres, J. Carrillo, BP)

Assume

- $a \geq a_0 > 0$ and $b > 0$
- the initial data is concentrated enough around $v = V_F$.

Then,

- there are **NO** global weak solutions
- larger nonlinear diffusion does not help

Noisy LIF networks

Theorem (M. Cáceres, J. Carrillo, BP)

Assume

- $a \geq a_0 > 0$ and $b > 0$
- the initial data is concentrated enough around $v = V_F$.

Then,

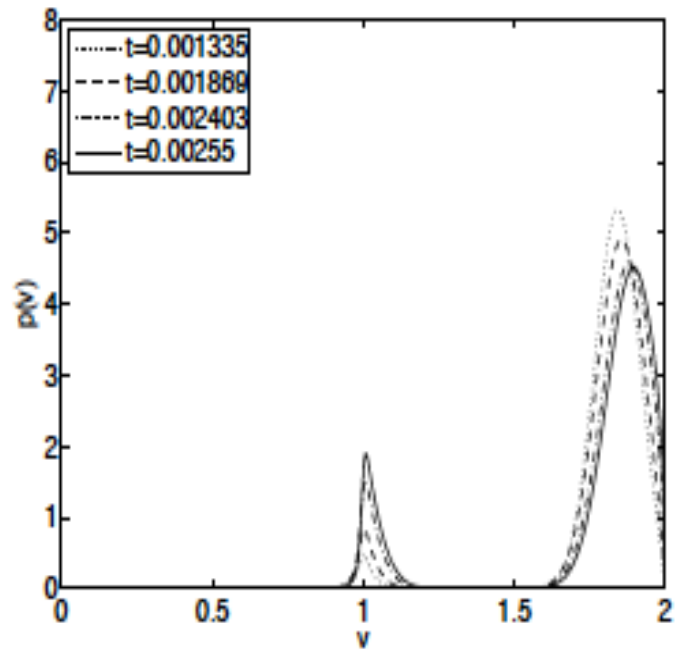
- there are **NO** global weak solutions
- larger nonlinear diffusion does not help

Possible interpretation

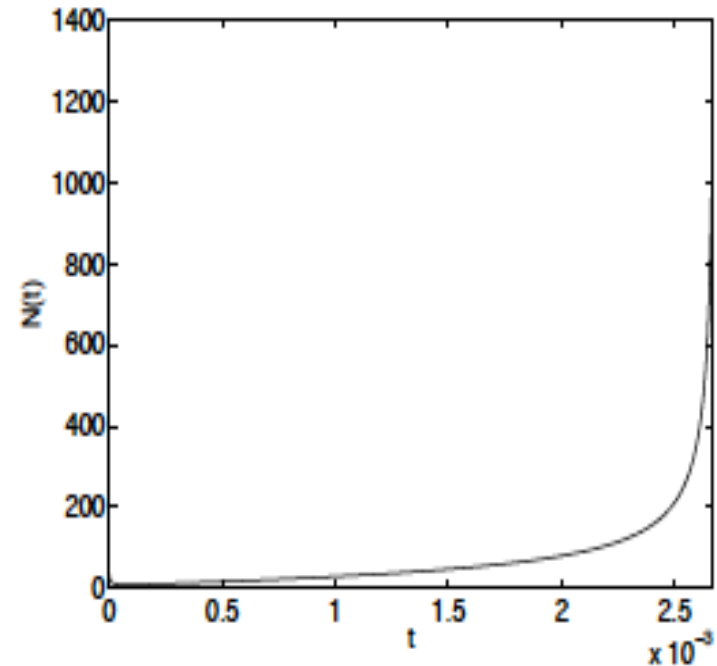
- $N(t) \rightarrow \rho \delta(t - t_{BU})$.
- partial synchronization (see S. Ha)

Noisy LIF networks

Numerical solution of the blow-up phenomena



probability density $n(v)$

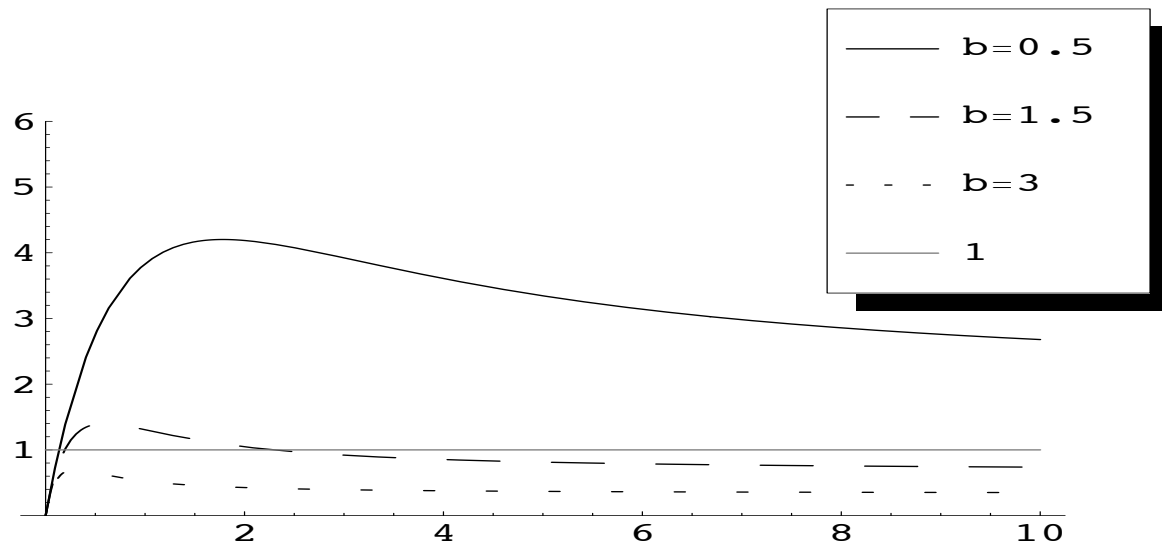


Total neuronal activity $N(t)$

Noisy LIF networks

Theorem (Steady states) For

- $b > 0$ small enough, there is a unique steady state
- $b > (V_F - V_R)$, $2ab < (V_F - V_R)^2 V_R$, then there are at least 2 steady states
- $b > 0$ large enough, there are no steady states.



Noisy LIF networks

Similarity with a Keller-Segel type model

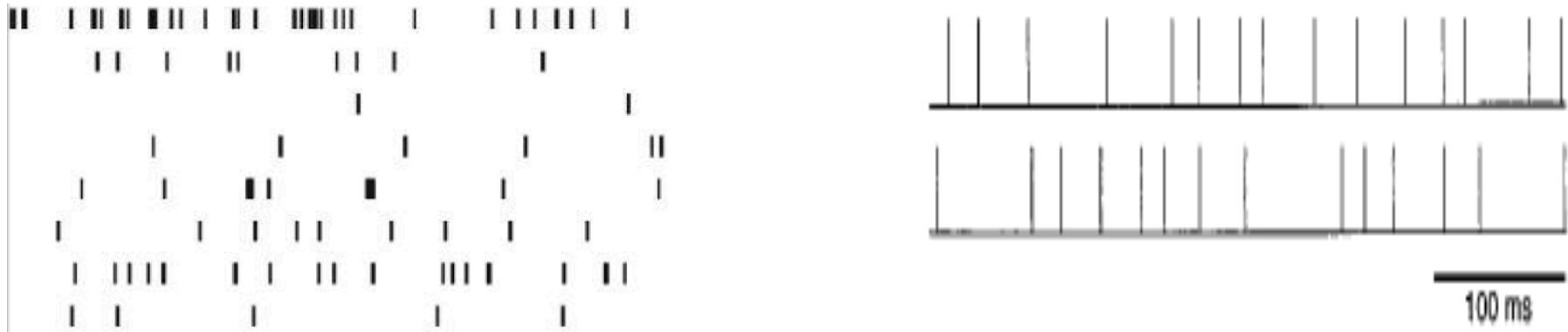
by V. Calvez and R. Voituriez

for microtubules arrangements on the membrane

$$\left\{ \begin{array}{l} \frac{\partial n(z,t)}{\partial t} - \frac{\partial}{\partial z} [\mu(t)n(z,t)] - \frac{\partial^2 n(z,t)}{\partial z^2} = 0, \quad z \geq 0, \\ \frac{\partial}{\partial z} n(0,t) + \mu(t)n(0,t) = 0, \\ \frac{d\mu(t)}{dt} = n(0,t) - \frac{\mu(t)}{L}. \end{array} \right.$$

- Blow-up for large mass
- Smooth solutions for small mass (and stable steady state)

Elapsed time structured model



K. Pakdaman, J. Champagnat, J.-F. Vibert have proposed to structure by time rather than potential which is a possible coding of neuronal information

Elapsed time structured model

- s represents the time elapsed since the last discharge
- $n(s, t)$ probability of finding a neuron in 'state' s at time t
- $p(s, N) \leq 1$ represents the firing rate of neurons in the 'state' s
- $N(t) = \text{activity of the network} + \text{external signaling}$

$$\left\{ \begin{array}{l} \frac{\partial n(s, t)}{\partial t} + \overbrace{\frac{\partial n(s, t)}{\partial s}}^{\text{elapsed time advances}} + \overbrace{p(s, bN(t)) n(s, t)}^{\text{firing neurons}} = 0, \\ n(s = 0, t) = \underbrace{\int_0^{+\infty} p(s, bN(t)) n(s, t) ds}_{\text{neurons reset}}, \end{array} \right.$$

This model always satisfies $\int_0^{+\infty} n(s, t) ds = 1.$

Elapsed time structured model

- s represents the time elapsed since the last discharge
- $n(s, t)$ probability of finding a neuron in 'state' s at time t
- $p(s, N) \leq 1$ represents the firing rate of neurons in the 'state' s '
- being given a total activity N

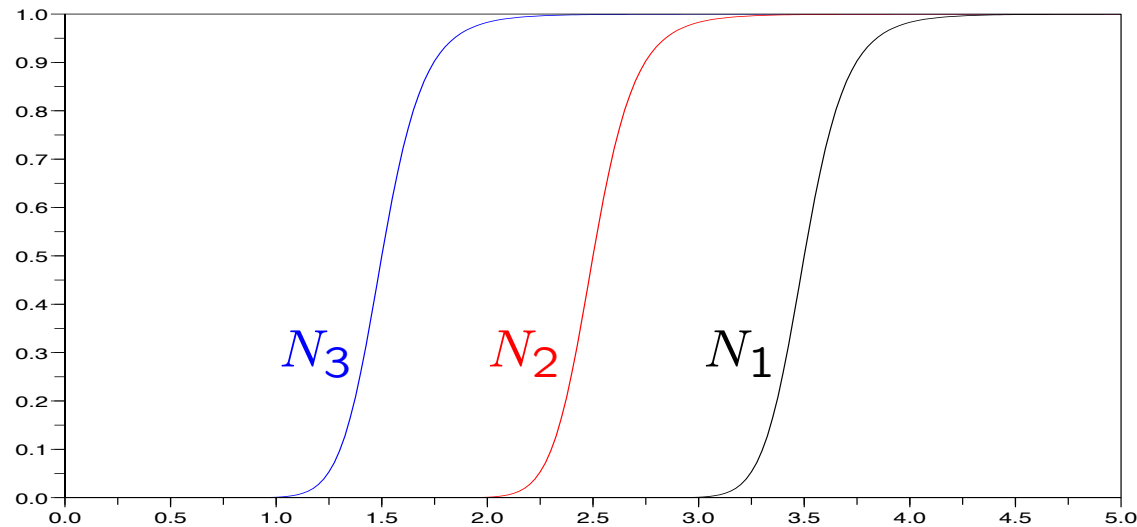
$$\begin{cases} \frac{\partial n(s, t)}{\partial t} + \frac{\partial n(s, t)}{\partial s} + p(s, bN(t)) n(s, t) = 0, \\ n(s = 0, t) = \int_0^{+\infty} p(s, bN(t)) n(s, t) ds, \end{cases}$$

$$N(t) := n(s = 0, t)$$

- $b > 0$ connectivity of the network
- excitatory neurons are represented by $\frac{\partial p(s, N)}{\partial N} > 0$

Elapsed time structured model

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + p(s, bN(t)) n(s,t) = 0, \\ n(s=0, t) = \int_0^{+\infty} p(s, bN(t)) n(s,t) ds, \end{cases}$$



the function $s \mapsto p(s, N)$ (refractory state+ fast transition)

Elapsed time structured model

With *synaptic integration*

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + p(s, X(t)) n(s,t) = 0, \\ N(t) := n(s=0, t) = \int_0^{+\infty} p(s, X(t)) n(s,t) ds, \end{cases}$$

$$X(t) := b \int_0^t N(t-u) \omega(u) du.$$

Elapsed time structured model

$$\begin{cases} \frac{\partial n(s,t)}{\partial t} + \frac{\partial n(s,t)}{\partial s} + p(s, bN(t)) n(s,t) = 0, \\ \frac{1}{b}N(t) := n(s=0, t) = \int_0^{+\infty} p(s, bN(t)) n(s,t) ds, \end{cases}$$

Properties

- $n \geq 0$, $\int_0^{\infty} n(s,t) ds = 1$,
- $N(t) \leq 1$, $n(s,t) \leq 1$,
- there is a unique solution,

Linear case For $p \equiv p(s)$ then

- $n(s,t) \xrightarrow[t \rightarrow \infty]{} P(s)$ the unique steady state.

Elapsed time structured model

The proof goes through Generalized Relative Entropy

$$\frac{d}{dt} \int_0^\infty \Phi(s) P(s) H\left(\frac{n(s,t)}{P(s)}\right) ds \leq 0,$$

for $H(\cdot)$ convex.

Elapsed time structured model

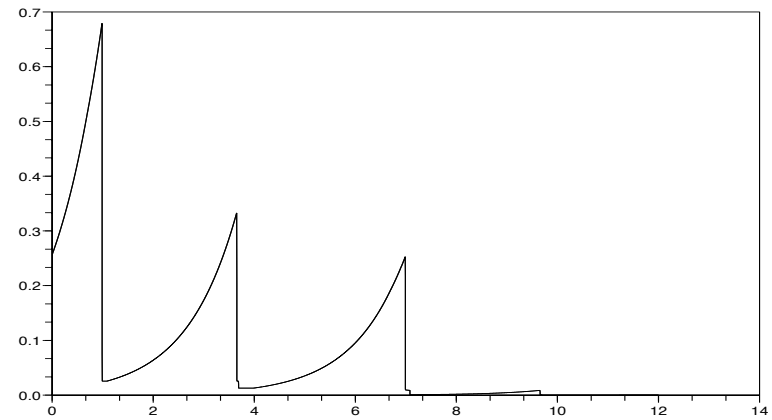
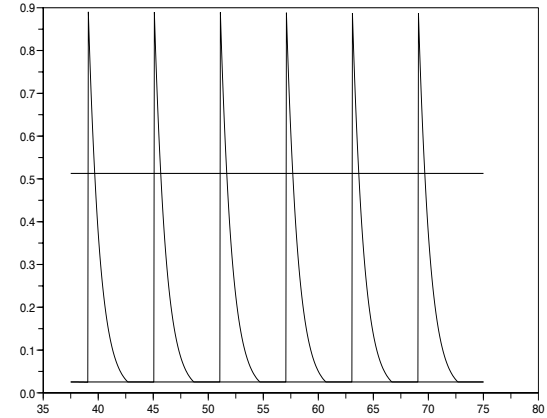
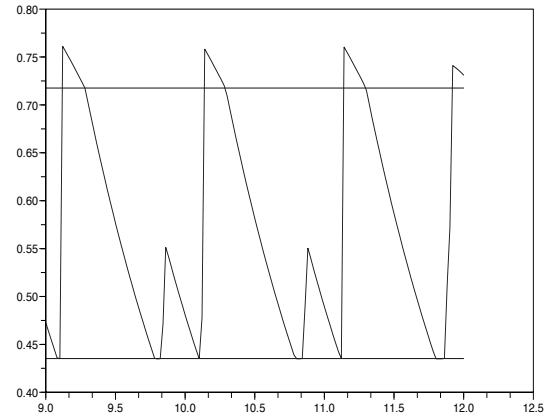
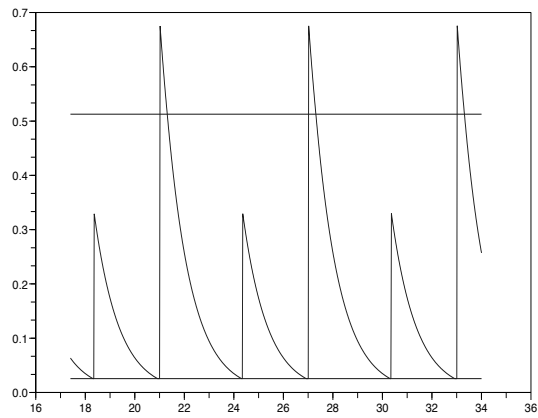
Properties

- For small or large connectivity (b small or large) then desynchronization still holds

$$n(s, t) \xrightarrow[t \rightarrow \infty]{} P_b(s)$$

- There are several periodic solutions (explicit),
- These are stable (observed numerically).

Elapsed time structured model



Conclusion

THANKS TO MY COLLABORATORS

M. J. Carceres (U.Granada), J. A. Carrillo (U. A. Barcelona)
(for I & F ; J. Mathematical Neurosciences 2011)

D. Smets (work in preparation)

K. Pakdaman, D. Salort (Inst. J. Monod, U. Paris Diderot)
(for elapsed time ; Nonlinearity 2010)

THANK YOU ALL