

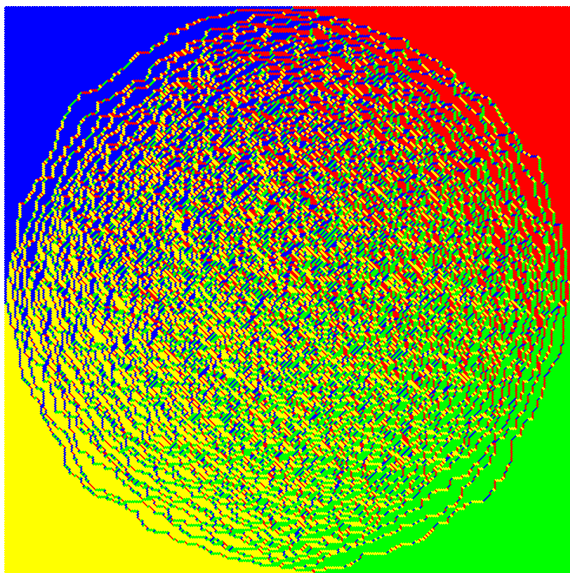
Limit shapes  
ICERM, April 13-17, 2015

# Limit shapes for interacting particle systems and their universal fluctuations

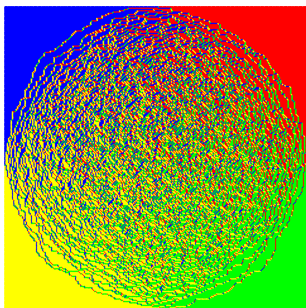
P.L. Ferrari in collaboration with A. Borodin



<http://wt.iam.uni-bonn.de/~ferrari>



An Aztec diamond of size  $N = 240$



The border of the four regular facets, as the size  $N \rightarrow \infty$ :

- has a circular limit shape (aka arctic circle)
 

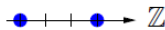
Jockush, Propp, Shor'98
- the fluctuations of border of the four facets are  $\mathcal{O}(N^{1/3})$  and (GUE) Tracy-Widom distributed
- As a process, it converges to the  $\text{Airy}_2$  process on the  $(N^{2/3}, N^{1/3})$  scale
 

Johansson'03

- TASEP: **Totally Asymmetric Simple Exclusion Process**

- **Configurations**

$$\eta = \{\eta_j\}_{j \in \mathbb{Z}}, \quad \eta_j = \begin{cases} 1, & \text{if } j \text{ is occupied,} \\ 0, & \text{if } j \text{ is empty.} \end{cases}$$



1 0 0 1  $\eta$

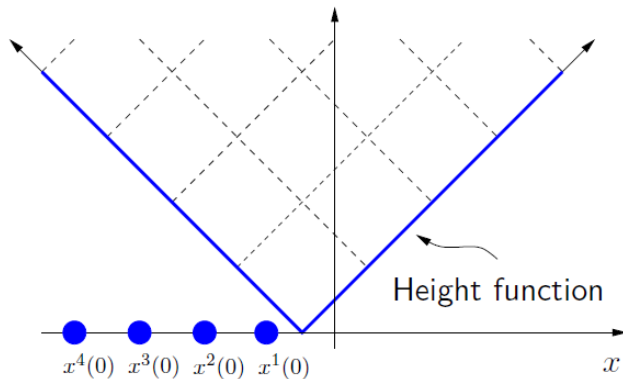
rate 1



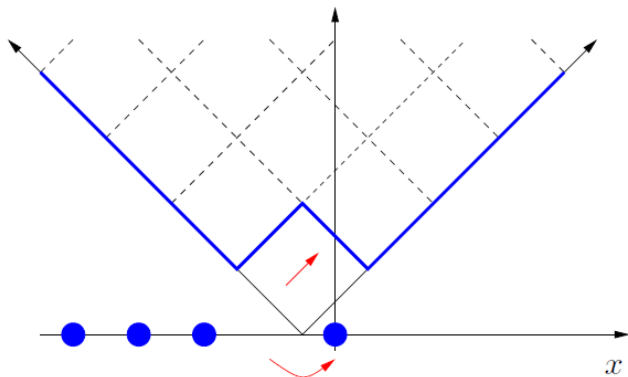
- **Dynamics**

Independently, particles jump on the right site with rate 1, provided the right is empty.

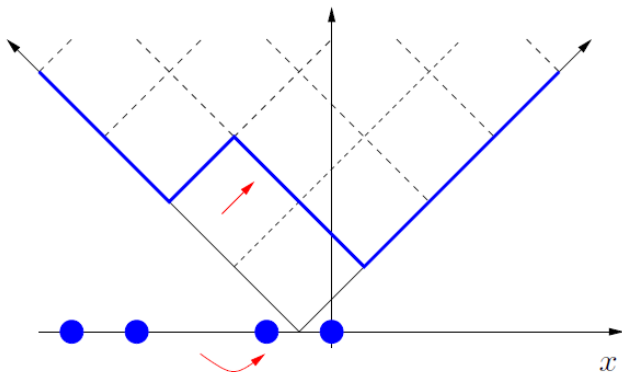
- ⇒ Particles are ordered: position of particle  $n$  is  $x^n(t)$
- Step initial condition is  $x^n(0) = -n, n \geq 1$ .



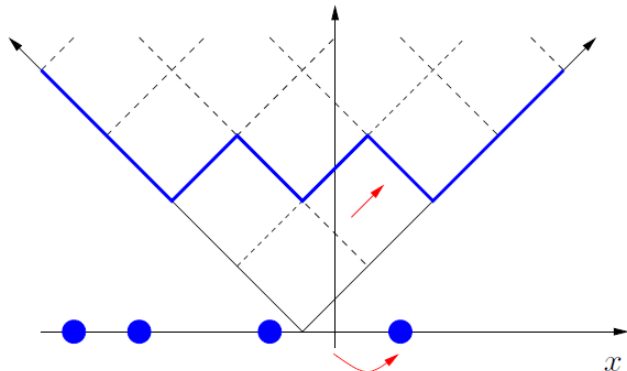
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Some known asymptotic results:

- law of large number:  $\lim_{t \rightarrow \infty} x^{\eta t}(t)/t = 1 - 2\sqrt{\eta}, \eta \in [-1, 1]$

Rost'81

- the fluctuations of particles are  $\mathcal{O}(t^{1/3})$  and (GUE)

Tracy-Widom distributed

- As a process in  $n$ , it converges to the  $\text{Airy}_2$  process on the  $(t^{2/3}, t^{1/3})$  scale

Johansson'03 (LPP); Borodin, Ferrari'07 (TASEP)

- Given a height function of a model in the Kardar-Parisi-Zhang universality class in one-dimension:  $x \mapsto h(x, t)$  (example:  $n \mapsto x^n(t)$ )

- Deterministic limit shape

$$h_{\text{ma}}(\xi) = \lim_{t \rightarrow \infty} h(\xi t, t)/t$$

- Stationary spatial diffusivity

$$A = \lim_{x \rightarrow \infty} \frac{\lim_{t \rightarrow \infty} \text{Var}(h(\xi t, t) - h(\xi t + x, t))}{|x|}$$

- Define further

$$\lambda = h''_{\text{ma}}(\xi) \quad \text{and} \quad \Gamma = |\lambda|A^2$$

- Rescaled process

$$h_t^{\text{resc}}(u) := \frac{h(\xi t + ut^{2/3}, t) - th_{\text{ma}}(\xi + ut^{-1/3})}{t^{1/3}}.$$

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$$h_t^{\text{resc}}(u) := \frac{h(\xi t + ut^{2/3}, t) - th_{\text{ma}}(\xi + ut^{-1/3})}{t^{1/3}}.$$

- If  $h''_{\text{ma}}(\xi) \neq 0$ , one expects the following:

$$\lim_{t \rightarrow \infty} h_t^{\text{resc}}(u) = \text{sgn}(\lambda)(\Gamma/2)^{1/3} \mathcal{A}_2 \left( \frac{Au}{2\Gamma^{2/3}} \right)$$

where  $\mathcal{A}_2$  is the Airy<sub>2</sub> process

Prähofer, Spohn'02

- For flat interfaces (i.e., if  $h''(\xi) = 0$ ) one has similar formulas but with either the Airy<sub>1</sub> process

Sasamoto'05; Borodin, Ferrari, Prähofer, Sasamoto'06

or the Airy<sub>stat</sub> depending on the initial conditions

Baik, Ferrari, Pécché'09

- With Alexei Borodin, in [Anisotropic growth of random surfaces in 2+1 dimensions](#) (arXiv:0804.3035), we introduced and studied a model of interacting particles in  $2 + 1$ -dimensions
- In discrete time, we have either *parallel update* or *sequential update*
- A discrete time *parallel update* includes (as different space-time projections) the Aztec diamond and the discrete time TASEP simultaneously

- The **state space** of our model is the Gelfand-Tsetlin pattern

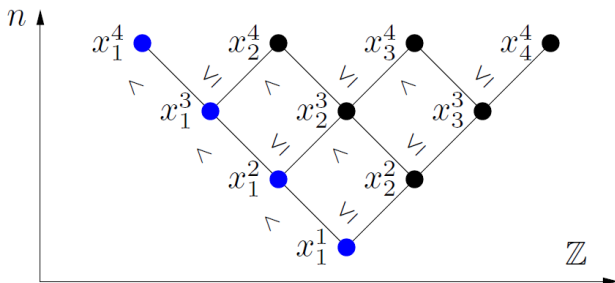
$$\text{GT}_N = \{X^N = (x^1, \dots, x^N); x^n = (x_1^n, \dots, x_n^n) \mid x^n \prec x^{n+1}, \forall n\}$$

where

$$x^n \prec x^{n+1} \Leftrightarrow x_1^{n+1} < x_1^n \leq x_2^{n+1} < x_2^n \leq \dots < x_n^n \leq x_{n+1}^{n+1}$$

means that  $x^n$  and  $x^{n+1}$  **interlace**.

- $x^n$  is called **configuration at level  $n$**



- The Markov chain at level  $n$  (discrete time) is given by  $x_1^n, \dots, x_n^n$  being one-sided random walk **conditioned to stay forever** in

$$W_n = \{x^n \in \mathbb{Z}^n \mid x_1^n < x_2^n < \dots < x_n^n\}.$$

- It is the Doob h-transform of the free walk with  $h$  function the Vandermonde determinant

$$\Delta_n(x^n) = \prod_{1 \leq i < j \leq n} (x_j^n - x_i^n),$$

i.e., it has the one-time transition probability given by

$$P_n(x^n, y^n) = \frac{\Delta_n(y^n)}{\Delta_n(x^n)} \det(P(x_i^n, y_j^n))_{i,j=1}^n$$

with  $P(x, y) = p\delta_{y,x+1} + (1-p)\delta_{y,x}$ .

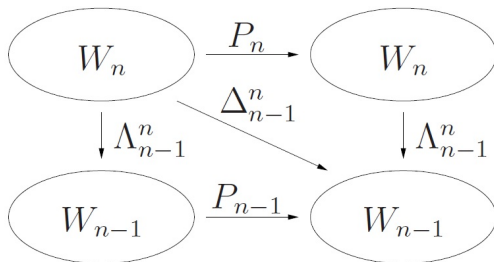
- The chain at fixed time  $t$  is the one that, given  $x^N$ , it generates the uniform measure on the interlacing configurations:

$$\begin{aligned}
 \Lambda_{n-1}^n(x^n, x^{n-1}) &:= \mathbb{P}(x^{n-1} | x^n) \\
 &= \frac{\#\text{GT}_{n-1} \text{ with given } x^{n-1}}{\#\text{GT}_n \text{ with given } x^n} \mathbb{1}_{x^{n-1} \prec x^n} \\
 &= (n-1)! \frac{\Delta_{n-1}(x^{n-1})}{\Delta_n(x^n)} \mathbb{1}_{x^{n-1} \prec x^n}
 \end{aligned}$$

- The key property used below is the **intertwining property** of the chains:

Diaconis, Fill '90

$$\Delta_{n-1}^n := P_n \Lambda_{n-1}^n = \Lambda_{n-1}^n P_{n-1}$$





- The sequential update is the following:
  - $x^1(t) \rightarrow x^1(t+1)$  according to  $P_1(x^1(t), x^1(t+1))$ ,
  - $x^2(t) \rightarrow x^2(t+1)$  to be the middle point of the chain  $(P_2 \circ \Lambda_1^2)(x^2(t), x^1(t+1))$
  - and so on

$$\begin{array}{ccc}
 x^3(t) & \xrightarrow{P_3} & x^3(t+1) \\
 \downarrow \Lambda_2^3 & & \downarrow \Lambda_2^3 \\
 x^2(t) & \xrightarrow{P_2} & x^2(t+1) \\
 \downarrow \Lambda_1^2 & & \downarrow \Lambda_1^2 \\
 x^1(t) & \xrightarrow{P_1} & x^1(t+1)
 \end{array}$$

- Projection on  $\{x_1^1, x_1^2, \dots, x_1^N\}$  is TASEP in discrete time with sequential update



- There is a **class of measure** which form is **invariant under  $P_\Lambda^N$** .  
Let  $\mu_N(x^N)$  be a probability measure on  $W_N$  and define

$$M_N(X^N) := \mu_N(x^N) \Lambda_{N-1}^N(x^N, x^{N-1}) \cdots \Lambda_1^2(x^2, x^1).$$

Then, applying  $t$  times  $P_\Lambda^N$  we have

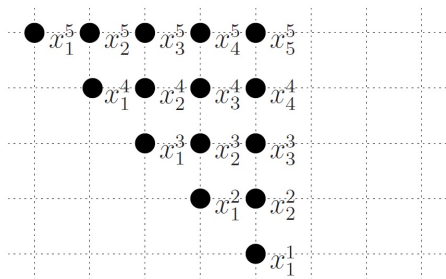
$$(M_N(P_\Lambda^N)^t)(Y^N) = (\mu_N(P_N)^t)(y^N) \Lambda_{N-1}^N(y^N, y^{N-1}) \cdots \Lambda_1^2(y^2, y^1)$$

- This is a consequence of the intertwining properties of the Markov chains!

- Consider further the "packed" initial condition:

$x_k^n(0) = -n + k$ ,  $1 \leq k \leq n \leq N$ . One can see that it can be written as above with  $\mu_N$  of the form

$$\mu_N(x^N) = \Delta_N(x^N) \det(\Psi_j(x_i^N, 0))_{i,j=1}^N.$$



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$x_k^n(0) = -n + k$ ,  $1 \leq k \leq n \leq N$ . One can see that it can be written as above with  $\mu_N$  of the form

$$\mu_N(x^N) = \Delta_N(x^N) \det(\Psi_j(x_i^N, 0))_{i,j=1}^N.$$

⇒ The measure at time  $t$  has the form

$$\prod_{n=1}^{N-1} \mathbb{1}_{[x^n < x^{n+1}]} \det(\Psi_j(x_i^N, t))_{i,j=1}^N$$

⇒ The measure at fixed level  $N$  and times  $t_1 < \dots < t_m$  has the form

$$\det(\Psi_j(x_i^N(t_1), t_1))_{i,j=1}^N \prod_{k=1}^{m-1} \det(P_{t_k, t_{k+1}}(x^N(t_k), x^N(t_{k+1})) \Delta_N(x^N(t_m)))$$

- Measure of this form have **determinantal correlations** as they are conditional  $L$ -ensembles

Borodin, Rains'06

Correlation structure of the **blue lozenges / particles**

Theorem (arXiv:0804.3035)

Consider any  $N$  triples  $(x_j, n_j, t_j)$  such that

$$t_1 \leq t_2 \leq \dots \leq t_N, \quad n_1 \geq n_2 \geq \dots \geq n_N.$$

Then,

$$\begin{aligned} \mathbb{P}(\text{at each } (x_j, n_j, t_j), j = 1, \dots, N, \\ \text{there exists a blue lozenge / particle}) \\ = \det[K(x_i, n_i, t_i; x_j, n_j, t_j)]_{1 \leq i, j \leq N} \end{aligned}$$

for an explicit kernel  $K$ .

Correlation structure of the three types of lozenges

Theorem (arXiv:0804.3035)

Consider any  $N$  triples  $(x_j, n_j, t_j)$  such that

$$t_1 \leq t_2 \leq \dots \leq t_N, \quad n_1 \geq n_2 \geq \dots \geq n_N.$$

Then,

$$\begin{aligned} \mathbb{P}(\text{at each } (x_j, n_j, t_j), j = 1, \dots, N, \\ \text{there exists a lozenge of color } c_j) \\ = \det[\tilde{K}(x_i, n_i, t_i, c_i; x_j, n_j, t_j, c_j)]_{1 \leq i, j \leq N} \end{aligned}$$

for an explicit kernel  $\tilde{K}$ .

- The parallel update is the following

$x^n(t) \rightarrow x^n(t+1)$  to be the middle point of the chain

$$(P_n \circ \Lambda_{n-1}^n)(x^n(t), x^{n-1}(t))$$


$$x^3(t) \xrightarrow{P_3} x^3(t+1)$$

$$\downarrow \Lambda_2^3$$

$$x^2(t) \xrightarrow{P_2} x^2(t+1)$$

$$\downarrow \Lambda_1^2$$

$$x^1(t) \xrightarrow{P_1} x^1(t+1)$$

- Projection on  $\{x_1^1, x_1^2, \dots, x_1^N\}$  is TASEP in discrete time with parallel update 

- This particle system is tightly related with the Aztec diamond:

- Start with packed initial condition:

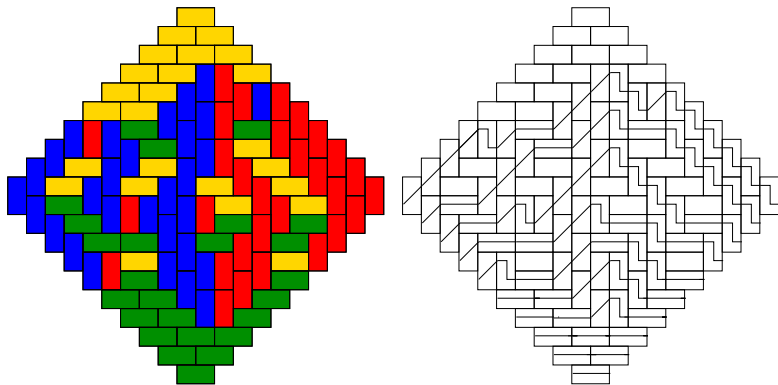
$$x^1 = (-1), \quad x^2 = (-2, -1), \quad x^3 = (-3, -2, -1).$$

- Extend our configuration space to:

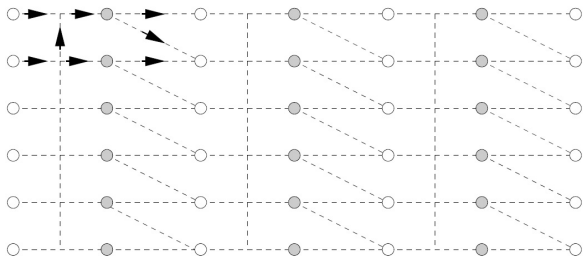
$$\begin{array}{ccccccc}
 x^3(t) & \xrightarrow{\text{Id}} & x^3(t+1) & \xrightarrow{\text{Id}} & & \xrightarrow{P_3} & \\
 \downarrow \Lambda_2^3 & & \downarrow \Lambda_2^3 & & \downarrow \Lambda_2^3 & & \downarrow \Lambda_2^3 \\
 y^3(t) & \xrightarrow{\text{Id}} & x^2(t) = y^3(t+1) & \xrightarrow{\text{Id}} & x^2(t+1) & \xrightarrow{P_2} & \\
 & & \downarrow \Lambda_1^2 & & \downarrow \Lambda_1^2 & & \downarrow \Lambda_1^2 \\
 & & y^2(t) & \xrightarrow{\text{Id}} & x^1(t) = y^2(t+1) & \xrightarrow{P_1} & x^1(t+1) \\
 & & & & \downarrow \Lambda_0^1 & & \downarrow \Lambda_0^1 \\
 & & & & y^1(t) = \emptyset & \xrightarrow{\text{Id}} & y^1(t+1) = \emptyset
 \end{array}$$



- Aztec diamond and line ensembles:

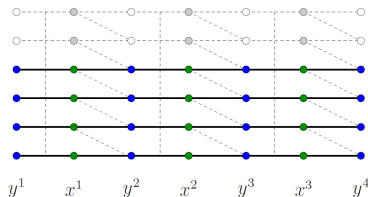


- A simple transformation of the Aztec line ensembles give a set of non-intersecting line ensembles on the following LGV graph with uniform weights



- Below we consider line ensembles with weight  $\alpha$  on vertical segments

- Possible configurations with their weights.



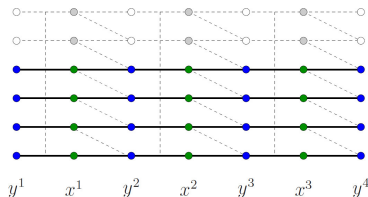
Time  $t = 0$

Weight 1

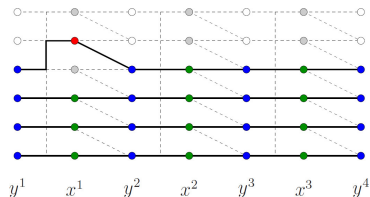
Probability 1

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Possible configurations with their weights.



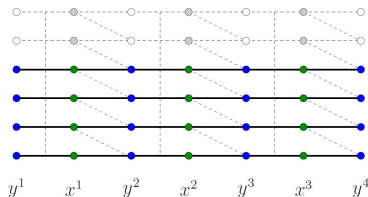
Time  $t = 0$   
 Weight 1  
 Probability 1



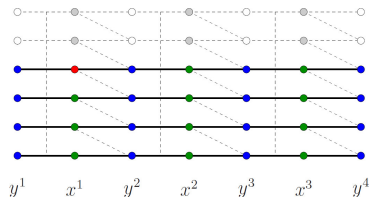
Time  $t = 1$   
 Weight  $\alpha$   
 Probability  $p$

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Possible configurations with their weights.



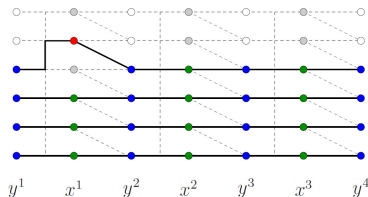
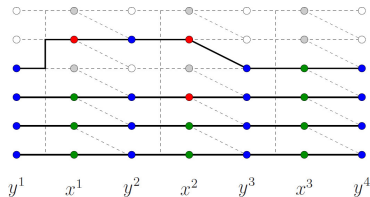
Time  $t = 0$   
 Weight 1  
 Probability 1



Time  $t = 1$   
 Weight 1  
 Probability  $1 - p$

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

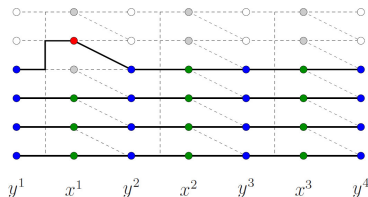
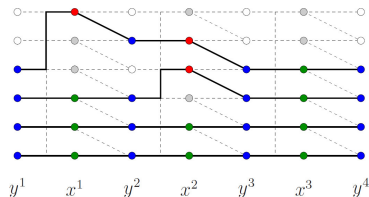
- Possible configurations with their weights.

Time  $t = 1$ Weight  $\alpha$ Probability  $p$ Time  $t = 2$ Weight  $\alpha$ Probability  $p \cdot (1 - p)^2$ 

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.



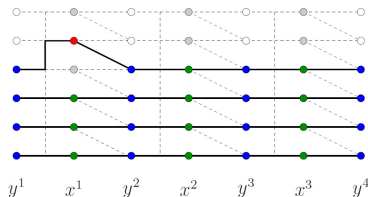
- Possible configurations with their weights.

Time  $t = 1$ Weight  $\alpha$ Probability  $p$ Time  $t = 2$ Weight  $\alpha^3$ Probability  $p \cdot p^2$ 

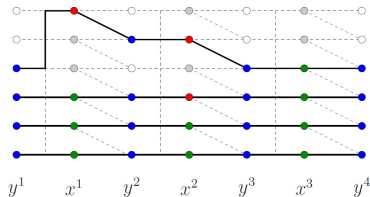
- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.



- Possible configurations with their weights.



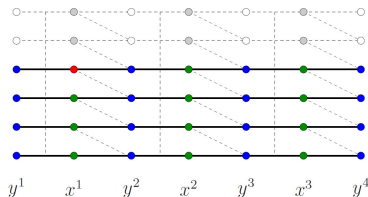
Time  $t = 1$   
 Weight  $\alpha$   
 Probability  $p$



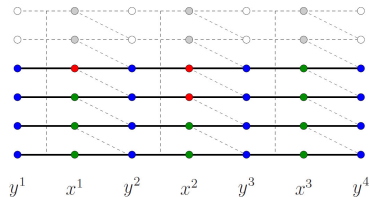
Time  $t = 2$   
 Weight  $\alpha^2$   
 Probability  $p \cdot p(1 - p)$

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Possible configurations with their weights.

Time  $t = 1$ 

Weight 1

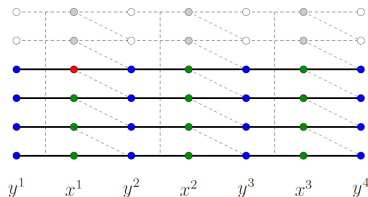
Probability  $1 - p$ Time  $t = 2$ 

Weight 1

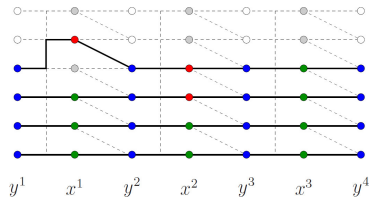
Probability  $(1 - p) \cdot (1 - p)^2$ 

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Possible configurations with their weights.

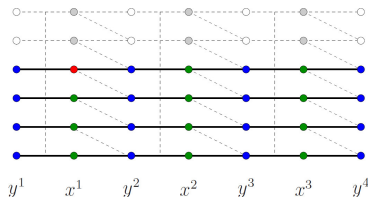
Time  $t = 1$ 

Weight 1

Probability  $1 - p$ Time  $t = 2$ Weight  $\alpha$ Probability  $(1 - p) \cdot p(1 - p)$ 

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

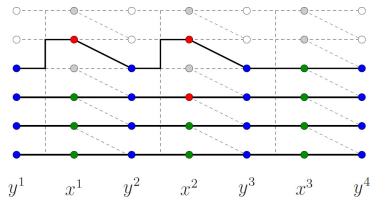
- Possible configurations with their weights.



Time  $t = 1$

Weight 1

Probability  $1 - p$



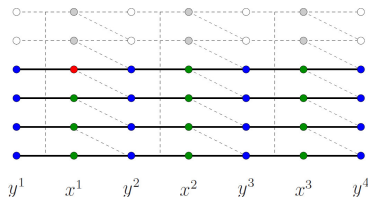
Time  $t = 2$

Weight  $\alpha^2$

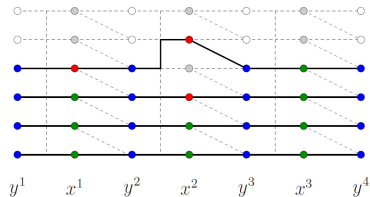
Probability  $(1 - p) \cdot p^2$

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Possible configurations with their weights.

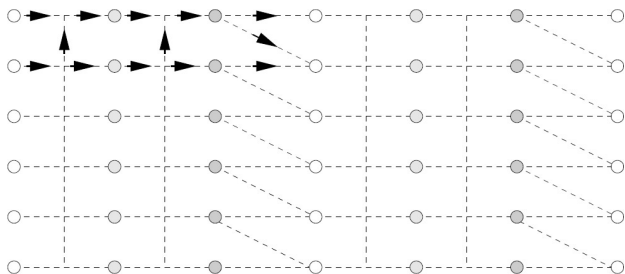
Time  $t = 1$ 

Weight 1

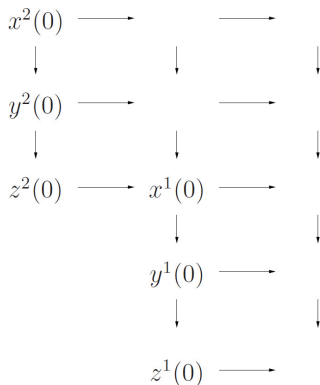
Probability  $1 - p$ Time  $t = 2$ Weight  $\alpha$ Probability  $(1 - p) \cdot (1 - p)p$ 

- The red points are the only that are not fixed by the boundary conditions: they form the particle process described above.

- Consider a simple generalization of the line ensembles by staying this time to this LGV graph



- The previous example fits in a dynamics with the following scheme



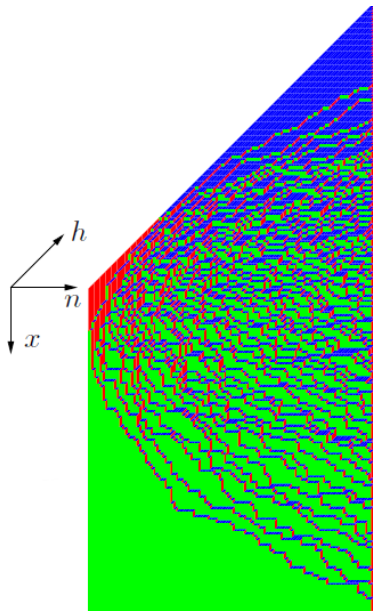
Macroscopic parametrization:

- $x = [-\eta L + \nu L]$
- $n = [\eta L]$
- $t = \tau L$

for a  $L \gg 1$ .

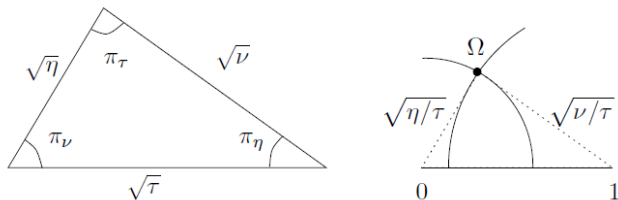
Asymptotic domain with “irregular”  
tiling (bordered by facets)

$$\mathcal{D} = \{(\nu, \eta, \tau), |\sqrt{\tau} - \sqrt{\eta}| < \sqrt{\nu} < \sqrt{\tau} + \sqrt{\eta}\}$$





Bulk:  $\mathcal{D} = \{(\nu, \eta, \tau) \in \mathbb{R}_+^3, |\sqrt{\tau} - \sqrt{\eta}| < \sqrt{\nu} < \sqrt{\tau} + \sqrt{\eta}\}$



Map  $\Omega : \mathcal{D} \rightarrow \mathbb{H} = \{z \in \mathbb{C} | \text{Im}(z) > 0\}$

Kenyon'04

- $\Omega$  is the critical point in the step descent analysis of the correlation kernel!
- $(\pi_\nu/\pi, \pi_\eta/\pi, \pi_\tau/\pi)$  are the **frequencies** of the three types of lozenge tilings: **Blue** for  $\pi_\eta$ , **Red** for  $\pi_\tau$ , **Green** for  $\pi_\nu$

- Limit shape:

$$\bar{h}(\nu, \eta, \tau) := \lim_{L \rightarrow \infty} \frac{\mathbb{E}(h((\nu - \eta)L, \eta L, \tau L))}{L} = \int_{\nu}^{(\sqrt{\tau} + \sqrt{\nu})^2} \frac{\pi_{\eta}(\nu', \eta, \tau)}{\pi} d\nu'$$

- The slopes are

$$\frac{\partial \bar{h}}{\partial \nu} = -\frac{\pi_{\eta}}{\pi}, \quad \frac{\partial \bar{h}}{\partial \eta} = 1 - \frac{\pi_{\nu}}{\pi}$$

- Growth velocity:

$$\frac{\partial \bar{h}}{\partial \tau} = \frac{\sin(\pi_{\nu}) \sin(\pi_{\eta})}{\pi \sin(\pi_{\tau})} = \frac{\text{Im}(\Omega)}{\pi}$$

## Theorem (arXiv:0804.3035)

For all  $(\nu, \eta, \tau) \in \mathcal{D}$ , denote  $\kappa = (\nu - \eta, \eta, \tau)$ . We have *moment convergence* of

$$\lim_{L \rightarrow \infty} \frac{h(\kappa L) - \mathbb{E}(h(\kappa L))}{\sqrt{c \ln L}} = \xi \sim \mathcal{N}(0, 1)$$

with  $c = 1/(2\pi^2)$  is *independent* of the macroscopic position in  $\mathcal{D}$ .

## Theorem (arXiv:0804.3035)

Consider any (disjoint)  $N$  triples  $\kappa_j = (\nu_j - \eta_j, \eta_j, \tau_j)$ , with  $(\nu_j, \eta_j, \tau_j) \in \mathcal{D}$ ,

$$\tau_1 \leq \tau_2 \leq \dots \leq \tau_N \quad \eta_1 \geq \eta_2 \geq \dots \geq \eta_N.$$

Set  $H_L(\kappa) := \sqrt{\pi} (h(\kappa L) - \mathbb{E}(h(\kappa L)))$ . Then,

$$\begin{aligned} & \lim_{L \rightarrow \infty} \mathbb{E}(H_L(\kappa_1) \cdots H_L(\kappa_N)) \\ &= \begin{cases} 0, & \text{odd } N, \\ \sum_{\text{pairings } \sigma} \prod_{j=1}^{N/2} G(\Omega_{\sigma(2j-1)}, \Omega_{\sigma(2j)}), & \text{even } N, \end{cases} \end{aligned}$$

with  $G(z, w) = -(2\pi)^{-1} \ln |(z - w)/(z - \bar{w})|$  is the Green function of the Laplacian on  $\mathbb{H}$  with Dirichlet boundary conditions.