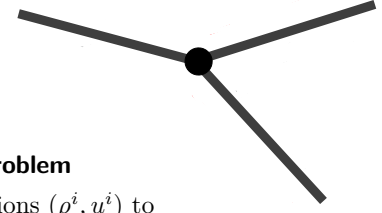


Entropy Methods for Gas Dynamics on Networks

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Joint work with Michael Herty, Michael Westdickenberg



Isentropic gas dynamics

We consider weak entropy solutions to

$$\begin{cases} \partial_t \rho^i + \partial_x(\rho u)^i & = 0 \\ \partial_t(\rho u)^i + \partial_x(\rho u^2 + \kappa \rho^\gamma)^i & = 0 \end{cases} \quad t, x > 0,$$

with density $\rho^i \geq 0$, flow velocity $u \in \mathbb{R}$, $\kappa > 0$, $1 < \gamma < 3$ and $i = 1, \dots, n$.

A kinetic BGK model [Bouchut, 1999]

$$\partial_t f^i + \xi \partial_x f^i = \frac{1}{\epsilon} (M[f^i] - f^i), \quad t, x > 0,$$

for $f^i = f^i(t, x, \xi) \in \mathbb{R}^2$, $f_0^i \geq 0$ and **Maxwellian** $M[f] = M(\rho_f, u_f, \xi)$ with moments $(\rho_f, \rho_f u_f) = \int_{\mathbb{R}} f(\xi) d\xi$.

Maximum energy dissipation problem

We select a coupling condition

$$f^i(t, x = 0, \xi) = \Psi^i[f^j(t, x = 0, \mathbb{R}^-), 1 \leq j \leq n](\xi)$$

for a.e. $t > 0, \xi > 0$, by solving the maximum energy dissipation problem

$$\begin{aligned} \inf \sum_{i=1}^d A^i \int_0^\infty \xi H(\Psi^i(\xi), \xi) d\xi \\ \text{s.t. } \sum_{i=1}^d A^i \left[\int_0^\infty \xi \Psi_0^i(\xi) d\xi + \int_{-\infty}^0 \xi g_0^i(\xi) d\xi \right] = 0 \end{aligned}$$

for given data $g^i(\xi)$, $\xi < 0$, where H denotes the kinetic energy to the BGK model. The solution is

$$\Psi^i(\xi) = M(\rho_*, u_* = 0, \xi) \quad \text{for } \xi > 0.$$

for suitable $\rho_* \geq 0$. Solutions to the kinetic BGK model on networks exist if initial total mass and energy are finite [2].

Interior relaxation

By **compensated compactness**, we obtain strong convergence towards weak entropy solutions to the system of isentropic gas dynamics [2].

A new coupling condition

The kinetic approach and assuming $\rho_*^\epsilon \rightarrow \rho_*$ in L^1_{loc} leads to the following definition of the generalized Riemann problem. Solutions exist and are unique for all initial data [3], not necessarily being subsonic.

References

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Generalized Riemann problem

1. Self-similar Lax solutions (ρ^i, u^i) to

$$(\rho^i, \rho^i u^i)(t = 0, x) = \begin{cases} (\rho_*, 0) & \text{if } x < 0, \\ (\rho_0^i, \rho_0^i u_0^i) & \text{if } x > 0. \end{cases}$$

2. The artificial density $\rho_* \geq 0$ is chosen such that

$$\sum_{i=1}^d A^i (\overline{\rho u})^i(t, 0+) = 0.$$

3. We restrict $(\rho^i, \rho^i u^i)$ to $x > 0$.

Comparison of coupling conditions

Conservation of total mass

$$\sum_{i=1}^d A^i (\rho u)^i(t, 0+) = 0.$$

Equal pressure [Banda, Herty, Klar, 2006]:

$$(\rho^\gamma)^i(t, 0+) = \mathcal{H}(t) \quad \text{for all } i.$$

Equal momentum flux [Colombo, Garavello, 2006]:

$$(\rho u^2 + \kappa \rho^\gamma)^i(t, 0+) = \mathcal{H}(t) \quad \text{for all } i.$$

Equal stagnation enthalpy [Reigstad, 2015]:

$$\left(\frac{u^2}{2} + \frac{\gamma \kappa}{\gamma - 1} \rho^{\gamma-1} \right)^i(t, 0+) = \mathcal{H}(t) \quad \text{for all } i.$$

A numerical example

Set $n = 3$, $\gamma = 2$, $\kappa = 5$, $A^i = 1$ with initial data

$$\begin{cases} \rho_{0,i} = 1 & i = 1, 2, 3, \\ u_{0,1} = -1, & u_{0,2} = u_{0,3} = 1/2 \end{cases}$$

The new coupling condition leads to physically meaningful waves and dissipates energy at the junction:

| pipe | new condition | pressure | momentum flux | stagnation enthalpy |
|--------------------|---------------|----------|---------------|---------------------|
| 1 | S | NW | R | R |
| 2 | R | NW | S | S |
| 3 | R | NW | S | S |
| energy dissipation | + | 0 | +/- | 0 |

NW=no wave, S=shock, R=rarefaction wave